

# Mathematical Analysis of High Aspect-Ratio Aircraft Wing in Subsonic Air Flow: Incompressible and Compressible Cases

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# **Flutter Analysis**

### Flutter

## Flutter Analysis

Definition

F-15B/FTF - II

Experiment

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Aeroelastic Problem

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Spectral Asymptotics

Analytic Mechanism

Possio Int. Eq.

- Flutter is an instability endemic to aircraft that occurs at high enough airspeed in subsonic flight and sets a "flutter boundary" on attainable airspeed in the subsonic regime.
  - The determination of aircraft flutter characteristics is one of the most important safety aspects in the aircraft design and analysis process. Damage inflicted by flutter is extremely costly.
- Flutter analysis is ranked by NASA aeronautics as one of the top research projects
- "Some fear flutter because they do not understand it" said the famous aerodynamist Theodore von Karman. "And some fear it," he added "because they do."



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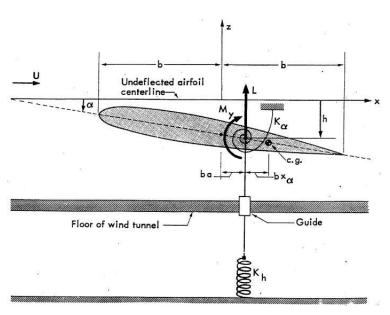
Possio Int. Eq.

"Flutter is the onset, beyond some speed -altitude combinations, of unstable and destructive vibrations of a lifting surface in an airstream. Flutter most commonly encountered on bodies subjected to large lateral aerodynamic loads of lift type, such as wings, tails, and control surfaces"

# $\dot{\mathbf{\Psi}} = \mathbf{i}\mathcal{L}\mathbf{\Psi} + \int_{0}^{\mathbf{t}} \mathcal{F}(\mathbf{t} - \sigma)\dot{\mathbf{\Psi}}(\sigma)d\sigma$

Evolution-convolution equation

# Theodorsen - Garrick experiment





# F-15B/FTF - II in Flight

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NASA Dryden Flight Research Center Photo Collection http://www.dfrc.nasa.gov/Gallery/Photo/index.html NASA Photo: EC05–0028–18 Date: February 14, 2005 Photo By: Carla Thomas

All six divots of thermal insulation foam have been ejected from the flight test fixture on NASA's F-15B testbed as it returns from a LIFT experiment flight.



# Flight Experiment Setup

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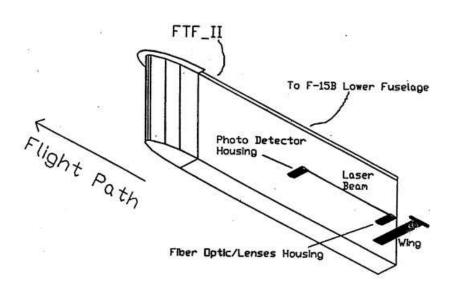
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# • Two Experiments conducted Simultaneously:

- Aeroelasticity Experiment with Flexible Wing
- Testing of New Gust Monitoring Device





# Some Facts about the flights

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- Taking into account all four flights, the experiment flew almost 7 hours.
- A total of 3.4 billion data points were sampled and recorded during the four flights by the high speed on-board recorder alone.
- Mach=0.8 @ 2000 feet were the maximum speed and altitudes achieved.
- Generated high angle of attack data & unique dynamic stall data ⇒ sparked new research interest.



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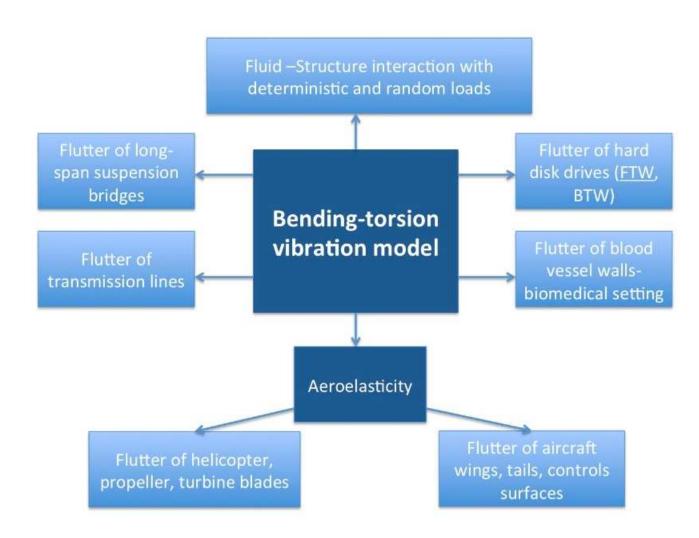
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# Models to be discussed

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Model of high aspect-ratio aircraft wing in subsonic, inviscid, incompressible air flow

Same type wing in subsonic, inviscid, compressible air flow.
 Aerodynamic loads and Possio integral equation



# Some of the main results to be presented

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Explicit asymptotic formulas for ground vibrations modes and corresponding eigenfunctions

 Explicit asymptotic formulas for aeroelastic modes and corresponding mode shapes

Analytic mechanism generating fluttering modes ("flutter matrix")

 Practical efficiency of explicit asymptotic formulas for aeroelastic modes; comparison with numerical results



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# **Aeroelastic Problem**

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bending torsion angle 
$$\mathbf{X}(\mathbf{x},\mathbf{t}) = (h(x,t), \ \alpha(x,t))^{\mathbf{T}}, \qquad -\mathbf{L} \leq \mathbf{x} \leq \mathbf{0}, \ \mathbf{t} \geq \mathbf{0},$$

Model: system of two coupled damped integro-differential equations

$$\begin{split} &(\mathbf{M_s} - \mathbf{M_a}) \ddot{\mathbf{X}}(\mathbf{x}, \mathbf{t}) + (\mathbf{D_s} - \mathbf{u} \mathbf{D_a}) \dot{\mathbf{X}}(\mathbf{x}, \mathbf{t}) + \\ &(\mathbf{K_s} - \mathbf{u^2} \mathbf{K_a}) \mathbf{X} = \left[\mathbf{f_1}(\mathbf{x}, \mathbf{t}), \mathbf{f_2}(\mathbf{x}, \mathbf{t})\right]^\mathbf{T}. \end{split}$$

where  $\mathbf{u} > 0$  - stream velocity

$$\mathbf{M_s} = \left[ egin{array}{ccc} \mathbf{m} & \mathbf{S} \\ \mathbf{S} & \mathbf{I} \end{array} 
ight], \qquad \mathbf{M_a} = (-\pi 
ho) \left[ egin{array}{ccc} \mathbf{1} & -\mathbf{a} \\ -\mathbf{a} & (\mathbf{a^2} + \mathbf{1/8}) \end{array} 
ight],$$

 ${\bf m}$  - density of the structure,  $~{\bf S}$  - mass moment,  ${\bf I}$  - moment of inertia;  $\rho\ll 1$  , a  $\in [-1,1]$  ;



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$$\mathbf{D_s} = \left[ egin{array}{cc} \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} \end{array} 
ight],$$

$$\mathbf{D_a} = (-\pi 
ho) \left[ egin{array}{cc} \mathbf{0} & \mathbf{1} \ -\mathbf{1} & \mathbf{0} \end{array} 
ight],$$

$$\mathbf{K_s} = \begin{bmatrix} \mathbf{E} \frac{\partial^4}{\partial \mathbf{x}^4} & \mathbf{0} \\ \mathbf{0} & -\mathbf{G} \frac{\partial^2}{\partial \mathbf{x}^2} \end{bmatrix}, \qquad \mathbf{K_a} = (-\pi \rho) \begin{bmatrix} 0 & 0 \\ 0 & -1 \end{bmatrix},$$

E - bending stiffness; G - torsion stiffness



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$$\begin{split} \mathbf{f_1}(\mathbf{x},\mathbf{t}) &= -2\pi\rho\int\limits_0^\mathbf{t} \left[\mathbf{u}\mathbf{C_2}(\mathbf{t}-\sigma) - \dot{\mathbf{C}}_3(\mathbf{t}-\sigma)\right]\mathbf{g}(\mathbf{x},\sigma)\mathbf{d}\sigma, \\ \mathbf{f_2}(\mathbf{x},\mathbf{t}) &= -2\pi\rho\int\limits_0^\mathbf{t} \left[\mathbf{1}/2\mathbf{C_1}(\mathbf{t}-\sigma) - \mathbf{a}\mathbf{u}\mathbf{C_2}(\mathbf{t}-\sigma) + \mathbf{a}\dot{\mathbf{C}}_3(\mathbf{t}-\sigma) + \mathbf{u}\mathbf{C_3}(\mathbf{t}-\sigma) + \mathbf{u}\mathbf{C_4}(\mathbf{t}-\sigma) + \mathbf{1}/2\dot{\mathbf{C}}_5(\mathbf{t}-\sigma)\right]\mathbf{g}(\mathbf{x},\sigma)\mathbf{d}\sigma \end{split}$$

Aerodynamical functions:  $C_i$ ,  $i = 1 \dots 5$ ,

$$\mathbf{g}(\mathbf{x}, \sigma) = \mathbf{u} \, \dot{\alpha}(\mathbf{x}, \sigma) + \ddot{\mathbf{h}}(\mathbf{x}, \sigma) + \left(\frac{\mathbf{1}}{\mathbf{2}} - \mathbf{a}\right) \, \ddot{\alpha}(\mathbf{x}, \sigma)$$

# **Recursion Relations**

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$$\widehat{\mathbf{C}}_{\mathbf{1}}(\lambda) \quad = \quad \int\limits_{0}^{\infty} \mathbf{e}^{-\lambda \mathbf{t}} \mathbf{C}_{\mathbf{1}}(\mathbf{t}) \mathbf{dt} = \frac{\mathbf{u}}{\lambda} \; \frac{\mathbf{e}^{-\lambda/\mathbf{u}}}{\mathbf{K}_{\mathbf{0}}(\lambda/\mathbf{u}) + \mathbf{K}_{\mathbf{1}}(\lambda/\mathbf{u})}, \;\; \Re \lambda > \mathbf{0},$$

$$\mathbf{C_2}(\mathbf{t}) = \int_{0}^{\mathbf{t}} \mathbf{C_1}(\sigma) d\sigma,$$

$$\mathbf{C_3}(\mathbf{t}) = \int_{0}^{\mathbf{t}} \mathbf{C_1}(\mathbf{t} - \sigma)(\mathbf{u}\sigma - \sqrt{\mathbf{u^2}\sigma^2 + 2\mathbf{u}\sigma})d\sigma,$$

$$C_4(t) = C_2(t) + C_3(t),$$

$$\mathbf{C_5(t)} \quad = \quad \int\limits_0^t \mathbf{C_1(t-\sigma)}((1+u\sigma)\sqrt{u^2\sigma^2+2u\sigma}-(1+u\sigma)^2)d\sigma,$$

 $K_0(z)$  and  $K_1(z)$  - the modified Bessel functions (McDonald functions)



# **Boundary Conditions**

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## **Boundary Conditions**

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Possio Int. Eq.

At the end x = -L:

$$h(-L, t) = h'(-L, t) = \alpha(-L, t) = 0$$

At the end x = 0: self straining actuator action

$$\mathbf{E}\mathbf{h}''(\mathbf{0}, \mathbf{t}) + \beta \dot{\mathbf{h}}'(\mathbf{0}, \mathbf{t}) = \mathbf{0}, \quad \mathbf{h}'''(\mathbf{0}, \mathbf{t}) = \mathbf{0},$$

$$\mathbf{G}\alpha'(\mathbf{0}, \mathbf{t}) + \delta\dot{\alpha}(\mathbf{0}, \mathbf{t}) = \mathbf{0}, \quad \beta, \delta \in \mathfrak{C}^+ \cup \{\infty\},$$

If  $\beta>0$ , then  $\beta\equiv g_h$  - bending control gain, If  $\delta>0$ , then  $\delta\equiv g_\alpha$  - torsion control gain



# **Initial Conditions**

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(a)

**(b)** 

Possio Int. Eq.

$$\mathbf{h}(\mathbf{x}, \mathbf{0}) = \mathbf{h}_{\mathbf{0}}(\mathbf{x}), \quad \dot{\mathbf{h}}(\mathbf{x}, \mathbf{0}) = \mathbf{h}_{\mathbf{1}}(\mathbf{x}),$$

$$\alpha(\mathbf{x}, \mathbf{0}) = \alpha_{\mathbf{0}}(\mathbf{x}), \quad \dot{\alpha}(\mathbf{x}, \mathbf{0}) = \alpha_{\mathbf{1}}(\mathbf{x})$$

# **Assumptions**

$$\det \left[ egin{array}{cc} \mathbf{m} & \mathbf{S} \ \mathbf{S} & \mathbf{I} \end{array} 
ight] > 0,$$

$$0 < \mathrm{u} \leq rac{\sqrt{2\mathrm{G}}}{\mathrm{L}\sqrt{\pi
ho}} \quad ext{(divergence speed)}$$



# **Energy of vibration**

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$$\mathbf{E}(\mathbf{t}) = \frac{1}{2} \int_{-\mathbf{L}}^{\mathbf{0}} [\mathbf{E}|\mathbf{h}''(\mathbf{x},\mathbf{t})|^{2} + \mathbf{G}|\alpha'(\mathbf{x},\mathbf{t})|^{2} + \widetilde{\mathbf{m}}|\dot{\mathbf{h}}(\mathbf{x},\mathbf{t})|^{2} +$$

$$\tilde{\mathbf{I}}|\dot{\alpha}(\mathbf{x},\mathbf{t})|^2 + \tilde{\mathbf{S}}(\dot{\alpha}\dot{\overline{\mathbf{h}}} + \dot{\overline{\alpha}}\dot{\mathbf{h}}) - \pi\rho\mathbf{u}^2|\alpha(\mathbf{x},\mathbf{t})|^2]d\mathbf{x}$$

where 
$$\tilde{\mathbf{I}} = \mathbf{I} + \pi \rho (\mathbf{a^2} + 1/8), \quad \tilde{\mathbf{m}} = \mathbf{m} - \pi \rho, \quad \tilde{\mathbf{S}} = \mathbf{S} - \pi \rho \mathbf{a}$$

# **Energy dissipation**

$$\dot{\mathbf{E}}(\mathbf{t}) = -\mathbf{E}|\dot{\mathbf{h}}'(\mathbf{0}, \mathbf{t})|^{2} \Re \beta - \mathbf{G}|\dot{\alpha}(\mathbf{0}, \mathbf{t})|^{2} \Re \delta$$



# Initial - Boundary Value Problem

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$$\begin{cases} \tilde{\mathbf{m}}\ddot{\mathbf{h}} + \tilde{\mathbf{S}}\ddot{\alpha} + \mathbf{E}\mathbf{h}'''' + \pi\rho\mathbf{u}\dot{\alpha} &= \mathbf{f_1}(\mathbf{x}, \mathbf{t}), \quad -\mathbf{L} \leq \mathbf{x} \leq \mathbf{0}, \\ \tilde{\mathbf{S}}\ddot{\mathbf{h}} + \tilde{\mathbf{I}}\ddot{\alpha} - \mathbf{G}\alpha'' + \pi\rho\mathbf{u}^2\alpha - \pi\rho\mathbf{u}\mathbf{h} &= \mathbf{f_2}(\mathbf{x}, \mathbf{t}), \quad \mathbf{t} > \mathbf{0} \end{cases}$$

# **Boundary Conditions**

at 
$$\mathbf{x}=-\mathbf{L}$$
:  $\mathbf{h}(-\mathbf{L},\mathbf{t})=\mathbf{h}'(-\mathbf{L},\mathbf{t})=\alpha(-\mathbf{L},\mathbf{t})=\mathbf{0}$  at  $\mathbf{x}=\mathbf{0}$ :  $\mathbf{h}'''(\mathbf{0})=\mathbf{0}$ ; 
$$\begin{cases} \mathbf{E}\mathbf{h}''(\mathbf{0},\mathbf{t})+\beta\dot{\mathbf{h}}'(\mathbf{0},\mathbf{t})=\mathbf{0}\\ \mathbf{G}\alpha'(\mathbf{0},\mathbf{t})+\delta\dot{\alpha}(\mathbf{0},\mathbf{t})=\mathbf{0} \end{cases}$$

# **Initial Conditions**

$$\mathbf{h}(\mathbf{x}, \mathbf{0}) = \mathbf{h}_{\mathbf{0}}, \quad \dot{\mathbf{h}}(\mathbf{x}, \mathbf{0}) = \mathbf{h}_{\mathbf{1}}, \quad \alpha(\mathbf{x}, \mathbf{0}) = \alpha_{\mathbf{0}}, \quad \dot{\alpha}(\mathbf{x}, \mathbf{0}) = \alpha_{\mathbf{1}}$$

# **Operator Setting**

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# **State space** ↔ **Energy space**

 ${\cal H}$  - set of 4 component vector - valued functions

$$\Psi = (h, \dot{h}, \alpha, \dot{\alpha})^T \equiv (\psi_0, \psi_1, \psi_2, \psi_3)^T$$

satisfying

$$\psi_0(-L) = \psi_0'(-L) = \psi_2(-L) = 0$$

Norm of state space (Hilbert space):

$$||\Psi||_{\mathcal{H}}^2 = 1/2 \int\limits_{-L}^0 \left[ \mathbf{E} |\psi_0''(\mathbf{x})|^2 + \mathbf{G} |\psi_2'(\mathbf{x})|^2 + \widetilde{\mathbf{m}} |\psi_1(\mathbf{x})|^2 + \widetilde{\mathbf{I}} |\psi_3(\mathbf{x})|^2 \right]$$

$$+\tilde{\mathbf{S}}(\psi_{\mathbf{3}}(\mathbf{x})\bar{\psi}_{\mathbf{1}}(\mathbf{x}) + \bar{\psi}_{\mathbf{3}}(\mathbf{x})\psi_{\mathbf{1}}(\mathbf{x})) - \pi\rho\mathbf{u}^{2}|\psi_{\mathbf{2}}(\mathbf{x})|^{2}] d\mathbf{x}$$



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$$\mathbf{\dot{\Psi}} = \mathbf{i}\mathcal{L}\mathbf{\Psi} + \int_{0}^{\mathbf{t}} \mathcal{F}(\mathbf{t} - \sigma)\mathbf{\dot{\Psi}}(\sigma)\mathbf{d}\sigma$$

First order in time "evolution - convolution" equation in  ${\cal H}$ 

$$\dot{\Psi} = i\mathcal{L}_{\beta\delta}\Psi + \tilde{\mathcal{F}}\dot{\Psi}, \quad \Psi = (\psi_0, \psi_1, \psi_2, \psi_3)^T, \quad \Psi|_{t=0} = \Psi_0$$

 $\mathcal{L}_{eta\delta}$  - matrix differential operator in  $\mathcal{H}$ 

 $\widetilde{\mathcal{F}}$  - matrix integral operator in  $\mathcal{H}$ 



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# $\mathcal{L}_{eta\delta}$ - matrix differential operator in $\mathcal{H}$

$$\mathcal{L}_{eta\delta} = -\mathrm{i} \left[ egin{array}{cccc} 0 & 1 & 0 & 0 \ -rac{ ext{E} ilde{1}}{\Delta}rac{ ext{d}^4}{ ext{d} ext{x}^4} & -rac{\pi
ho ext{u} ilde{S}}{\Delta} & -rac{ ilde{S}}{\Delta}\left( ext{G}rac{ ext{d}^2}{ ext{d} ext{x}^2} + \pi
ho ext{u}^2
ight) & -rac{\pi
ho ext{u} ilde{I}}{\Delta} \ 0 & 0 & 0 & 1 \ rac{ ext{E} ilde{S}}{\Delta}rac{ ext{d}^4}{ ext{d} ext{x}^4} & rac{\pi
ho ext{u} ilde{m}}{\Delta} & rac{\widetilde{m}}{\Delta}\left( ext{G}rac{ ext{d}^2}{ ext{d} ext{x}^2} + \pi
ho ext{u}^2
ight) & rac{\pi
ho ext{u} ilde{S}}{\Delta} \end{array} 
ight]$$

# defined on the domain

$$\mathcal{D}(\mathcal{L}_{\beta\delta}) = \begin{cases} \mathbf{\Psi} \in \mathcal{H} : \ \psi_{\mathbf{0}} \in \mathbf{H^{4}}(-\mathbf{L}, \mathbf{0}), \ \psi_{\mathbf{1}} \in \mathbf{H^{2}}(-\mathbf{L}, \mathbf{0}), \\ \psi_{\mathbf{2}} \in \mathbf{H^{2}}(-\mathbf{L}, \mathbf{0}), \ \psi_{\mathbf{3}} \in \mathbf{H^{1}}(-\mathbf{L}, \mathbf{0}); \\ \psi_{\mathbf{1}}(-\mathbf{L}) = \psi'_{\mathbf{1}}(-\mathbf{L}) = \psi_{\mathbf{3}}(-\mathbf{L}) = \mathbf{0}; \ \psi''_{\mathbf{0}}(\mathbf{0}) = \mathbf{0}; \\ \mathbf{E}\psi''_{\mathbf{0}}(\mathbf{0}) + \beta\psi'_{\mathbf{1}}(\mathbf{0}) = \mathbf{0}, \ \mathbf{G}\psi'_{\mathbf{2}}(\mathbf{0}) + \delta\psi_{\mathbf{3}}(\mathbf{0}) = \mathbf{0} \end{cases}$$

$$\mathbf{\Delta} = \widetilde{\mathbf{m}}\mathbf{\tilde{I}} - \mathbf{\tilde{S}^2} > 0$$



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 $\widetilde{\mathcal{F}}$  - matrix integral operator in  $\mathcal{H}$ 

$$\widetilde{\mathcal{F}} = egin{bmatrix} 1 & 0 & 0 & 0 \ 0 & [ ilde{\mathbf{I}}( ilde{\mathbf{C}}_1*) - ilde{\mathbf{S}}( ilde{\mathbf{C}}_2*)] & 0 & 0 \ 0 & 0 & 1 & 0 \ 0 & 0 & 0 & [- ilde{\mathbf{S}}( ilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}( ilde{\mathbf{C}}_2*)] \end{bmatrix} imes$$

$$\left[egin{array}{cccc} 0 & 0 & 0 & 0 \ 0 & 1 & {f u} & (1/2-{f a}) \ 0 & 0 & 0 & 0 \ 0 & 1 & {f u} & (1/2-{f a}) \end{array}
ight]$$

Spectral properties of both the differential operator  $\mathcal{L}_{\beta\delta}$  and the integral operator  $\widetilde{\mathcal{F}}$  are of crucial importance for the representation of the solution



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$$\dot{\Psi} = i\mathcal{L}_{\beta\delta}\Psi + \widetilde{\mathcal{F}}\dot{\Psi}, \quad \Psi|_{t=0} = \Psi_0$$

Laplace transform representation for solution:

$$\widehat{\Psi}(\lambda) = \left(\lambda I - i\mathcal{L}_{\beta\delta} - \lambda \widehat{\mathcal{F}}(\lambda)\right)^{-1} (I - \widehat{\mathcal{F}}(\lambda))\Psi_0$$

Goal: find the solution in space - time domain, i.e., "calculate" the inverse Laplace transform of  $\widehat{\Psi}$ 

$$\mathcal{R}(\lambda) = \left(\lambda I - i\mathcal{L}_{\beta\delta} - \lambda \widehat{\mathcal{F}}(\lambda)\right)^{-1}$$

Generalized resolvent operator  $\Rightarrow$  analytic operator - valued function of  $\lambda$  on the complex plane having a branch - cut along the negative real semi - axis.

Poles of  $\mathcal{R}(\lambda)$  - discrete spectrum  $\leftrightarrow$  aeroelastic modes, Branch cut - continuous spectrum



# Spectral asymptotics of $\mathcal{L}_{\beta\delta}$

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Two - branch spectrum: If  $|\delta| 
eq \sqrt{\mathbf{G} \mathbf{ ilde{I}}}$ , then

 $\beta$ - branch

$$\mu_{\mathbf{n}}^{\beta} = (\mathbf{sgn}\,\mathbf{n})\pi^{2}/\mathbf{L}^{2}\sqrt{\mathbf{E}\tilde{\mathbf{I}}/\boldsymbol{\Delta}} (\mathbf{n} - \mathbf{1}/\mathbf{4})^{2} + \kappa_{\mathbf{n}}(\omega),$$
$$\omega = |\delta|^{-1} + |\beta|^{-1}, |\mathbf{n}| \to \infty,$$

where  $\mathbf{\Delta} = \widetilde{m} \mathbf{\tilde{I}} - \mathbf{\tilde{S}^2}$ , and

$$\sup_{\mathbf{n}\in\mathbb{Z}}\{|\kappa_{\mathbf{n}}(\omega)|\}=\mathbf{C}(\omega), \quad \mathbf{C}(\omega)\longrightarrow \mathbf{0} \text{ as } \omega\longrightarrow \mathbf{0}$$

 $\delta$ -branch

$$\mu_{\mathbf{n}}^{\delta} = \frac{\pi \mathbf{n}}{\mathbf{L}\sqrt{\tilde{\mathbf{I}}/\mathbf{G}}} + \frac{\mathbf{i}}{2\mathbf{L}\sqrt{\tilde{\mathbf{I}}/\mathbf{G}}} \, \ln \, \frac{\delta + \sqrt{\mathbf{G}\tilde{\mathbf{I}}}}{\delta - \sqrt{\mathbf{G}\tilde{\mathbf{I}}}} + \mathbf{O}(|\mathbf{n}|^{-1/2})$$



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Definition:  $\exists$  a nontrivial  $\Phi_n$  such that

$$\left(\lambda_n I - i\mathcal{L}_{\beta\delta} - \lambda_n \widehat{F}(\lambda_n)\right) \Phi_n = 0,$$

 $\Rightarrow \lambda_{\mathbf{n}}$ - aeroelastic mode,  $\Phi_{\mathbf{n}}$ - mode shape

# Theorem 1.

The set of aeroelastic modes  $\{\lambda_n\}$  is asymptotically close to the set  $\{i\mu_n\}$  where  $\{\mu_n\}$  are eigenvalues of  $\mathcal{L}_{\beta\delta}$ 

# Theorem 2.

- a)  $\mathcal{L}_{\beta\delta}$  is a closed linear operator with compact resolvent;
- b)  $\mathcal{L}_{\beta\delta}$  is nonselfadjoint unless  $\Re\beta=\Re\delta=0$ ;
- c) If  $\Re \beta > 0$  and  $\Re \delta > 0$ , then  $\mathcal{L}_{\beta \delta}$  dissipative:

$$\Im(\mathcal{L}_{\beta\delta}\Psi,\Psi)_{\mathcal{H}}\geq 0, \quad \Psi\in\mathcal{D}(\mathcal{L}_{\beta\delta});$$

d) Adjoint operator  $\mathcal{L}_{\beta\delta}^*$ ,  $\beta, \delta \mapsto -\bar{\beta}, -\bar{\delta}$ 



# Analytical mechanism generating fluttering modes

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# Structure of matrix integral operator

$$\mathfrak{S}(\lambda) = \lambda \mathbf{I} - \mathbf{i} \mathcal{L}_{\beta\delta} - \lambda \widehat{\mathcal{F}}(\lambda)$$

$$\widehat{\mathcal{F}} \Rightarrow egin{bmatrix} 1 & 0 & 0 & 0 & 0 \ 0 & [\widetilde{\mathbf{I}}(\widetilde{\mathbf{C}}1*) - \widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_2*)] & 0 & 0 \ 0 & 0 & 1 & 0 \ 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{m}}(\widetilde{\mathbf{C}}_2*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{0}}(\widetilde{\mathbf{C}}_1*) + [-\widetilde{\mathbf{S}}(\widetilde{\mathbf{C}}_1*) + \widetilde{\mathbf{C}}(\widetilde{\mathbf{C}}_1*)] \end{bmatrix} imes \ egin{bmatrix} \mathbf{0} & 0 & 0 & 0 & [-\widetilde{\mathbf{C}}(1*) + \widetilde{\mathbf{C}}(1*) + [-\widetilde{\mathbf{C}}(1*) + [-\widetilde{\mathbf{C}}$$



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 $\widehat{\mathcal{F}}(\lambda)$  is a Laplace transform of  $\mathcal{F}(\mathbf{t})\Rightarrow$  (kernel of convolution operator)

$$\widehat{\mathcal{F}}(\lambda) = egin{bmatrix} \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} \ \mathbf{0} & \mathcal{L}(\lambda) & \mathbf{u}\mathcal{L}(\lambda) & (\mathbf{1/2} - \mathbf{a})\mathcal{L}(\lambda) \ \mathbf{0} & \mathbf{0} & \mathbf{0} \ \mathbf{0} & \mathcal{N}(\lambda) & \mathbf{u}\mathcal{N}(\lambda) & (\mathbf{1/2} - \mathbf{a})\mathcal{N}(\lambda) \end{bmatrix}$$

$$\mathcal{L}(\lambda) = -\frac{2\pi\rho\mathbf{u}}{\lambda\mathbf{\Delta}} \left\{ -\tilde{\mathbf{S}}/\mathbf{2} + [\tilde{\mathbf{I}} + (\mathbf{1}/\mathbf{2} + \mathbf{a})\tilde{\mathbf{S}}]\mathbf{T}(\lambda/\mathbf{u}) \right\},\,$$

$$\mathcal{N}(\lambda) = -\frac{2\pi\rho\mathbf{u}}{\lambda\mathbf{\Delta}} \left\{ -\widetilde{\mathbf{m}}/\mathbf{2} + [\widetilde{\mathbf{S}} + (\mathbf{1}/\mathbf{2} + \mathbf{a})\widetilde{\mathbf{m}}]\mathbf{T}(\lambda/\mathbf{u}) \right\}$$



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# Theodorsen function:

$$\mathbf{T}(\mathbf{z}) = \frac{\mathbf{K_1}(\mathbf{z})}{\mathbf{K_0}(\mathbf{z}) + \mathbf{K_1}(\mathbf{z})}$$

 $K_0$ ,  $K_1$  the modified Bessel functions

$$\mathbf{K_n}(\mathbf{z}) = 1/2\pi \mathbf{i} \mathbf{e}^{\pi \mathbf{n} \mathbf{i}/2} \mathbf{H_n^{(1)}(iz)}$$

# **Definitions:**

$$\mathbf{K_0}(\mathbf{z}) = \sum_{\mathbf{m}=\mathbf{0}}^{\infty} \frac{\mathbf{z^{2m}}}{\mathbf{2^{2m}}(\mathbf{m}!)^2} \left( \psi(\mathbf{m}+\mathbf{1}) - \frac{1}{\mathbf{m}+\mathbf{1}} \ln(\mathbf{z}/\mathbf{2}) \right),$$

$$\mathbf{K_1}(\mathbf{z}) = \frac{1}{\mathbf{z}} + \frac{\mathbf{z}}{2} \sum_{\mathbf{m}=0}^{\infty} \frac{\mathbf{z^{2m}}}{\mathbf{2^{2m}}(\mathbf{m}!)^2(\mathbf{m}+1)} \left\{ \ln(\mathbf{z}/2) - \frac{1}{2} [\psi(\mathbf{m}+1) - \psi(\mathbf{m}+2)] \right\},$$

where  $\psi(\mathbf{m})$  is the digamma function



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Asymptotics as  $|\mathbf{z}| \to \infty$ 

$$T(z) = 1/2 + 1/16z + O(z^{-2})$$

$$\mathbf{V}(\mathbf{z}) = \mathbf{T}(\mathbf{z}) - \mathbf{1}/\mathbf{2} \to \mathbf{0} \quad \text{as } |\mathbf{z}| \to \infty$$

$$\mathfrak{S}(\lambda) = \lambda \mathbf{I} - \mathbf{i} \mathcal{L}_{\beta\delta} - \lambda \widehat{\mathcal{F}}(\lambda)$$



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$$\lambda \widehat{\mathcal{F}}(\lambda) = \mathfrak{M} + \mathfrak{N}(\lambda),$$

where

$$\mathfrak{M} = egin{bmatrix} 0 & 0 & 0 & 0 \ 0 & A & uA & (1/2-a)A \ 0 & 0 & 0 & 0 \ 0 & B & uB & (1/2-a)B \end{bmatrix}$$

$$\mathbf{A} = -\pi \rho \mathbf{u} / \mathbf{\Delta} [\tilde{\mathbf{I}} + (\mathbf{a} - 1/2)\tilde{\mathbf{S}}],$$
$$\mathbf{B} = \pi \rho \mathbf{u} / \mathbf{\Delta} [\tilde{\mathbf{S}} + (\mathbf{a} - 1/2)\tilde{\mathbf{m}}]$$



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$$\mathfrak{N}(\lambda) = egin{bmatrix} \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{A_1}(\lambda) & \mathbf{u}\mathbf{A_1}(\lambda) & (\mathbf{1/2-a})\mathbf{A_1}(\lambda) \\ \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{B_1}(\lambda) & \mathbf{u}\mathbf{B_1}(\lambda) & (\mathbf{1/2-a})\mathbf{B_1}(\lambda) \end{bmatrix}$$

$$\mathbf{A}_{1}(\lambda) = -2\pi \rho \mathbf{u} \Delta^{-1} \mathbf{V}(\lambda/\mathbf{u}) [\tilde{\mathbf{I}} + (\mathbf{a} + 1/2)\tilde{\mathbf{S}}],$$

$$\mathbf{B_1}(\lambda) = 2\pi \rho \mathbf{u} \Delta^{-1} \mathbf{V}(\lambda/\mathbf{u}) [\mathbf{\tilde{S}} + (\mathbf{a} + 1/2)\mathbf{\tilde{m}}]$$

$$\|\mathfrak{N}(\lambda)\| \le C|\lambda^{-1}|, \quad |\lambda| \to \infty$$



# Numerical results for the eigenvalues of the operator $i\mathcal{L}_{\beta\delta}+\mathfrak{M}$

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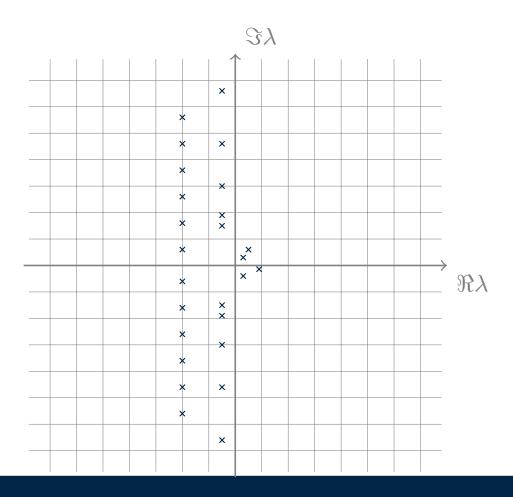
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Recall:  $\mathcal{R}(\lambda) = (\lambda I - i\mathcal{L}_{\beta\delta} - \mathfrak{M} - \mathfrak{N}(\lambda))^{-1}$  $\mathfrak{N}(\lambda)$ -Asymptotically small

M-"Flutter matrix"





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$$\mathfrak{S}(\lambda) = \lambda \mathbf{I} - \mathbf{i} \mathcal{L}_{\beta\delta} - \mathfrak{M} - \mathfrak{N}(\lambda), \quad \mathcal{R}(\lambda) = \mathfrak{S}(\lambda)^{-1}$$

Spectrum:  $\{\lambda_{\mathbf{n}}^{\beta}\}_{\mathbf{n}\in\mathbb{Z}}\cup\{\lambda_{\mathbf{n}}^{\delta}\}_{\mathbf{n}\in\mathbb{Z}}$  (aeroelastic modes)

Mode shapes:  $\{\Phi_{\mathbf{n}}^{\beta}\}_{\mathbf{n}\in\mathbb{Z}}\cup\{\Phi_{\mathbf{n}}^{\delta}\}_{\mathbf{n}\in\mathbb{Z}}$ 

# **Properties of mode shapes**

- 1. Minimality
- 2. Completeness in H
- 3. Riesz basis property



# **Contour Integration**

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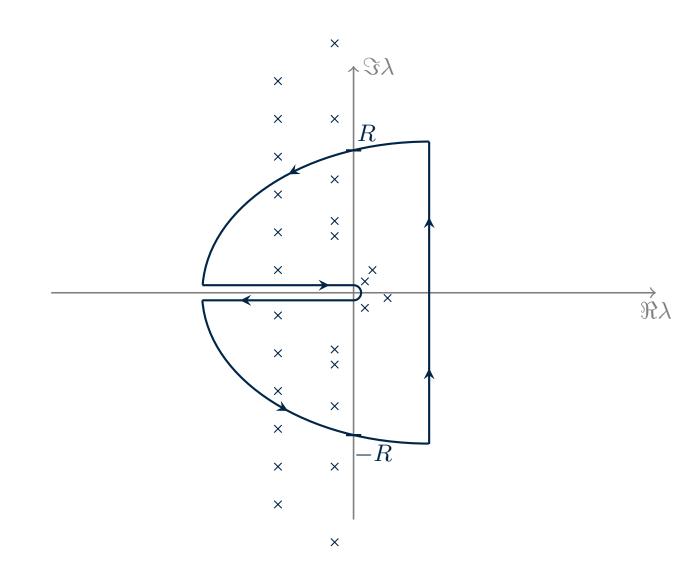
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$$\begin{split} \boldsymbol{\Psi}(\mathbf{x},\mathbf{t}) &= & \sum_{\mathbf{n} \in \mathbb{Z}'} \mathbf{e}^{\lambda_{\mathbf{n}}^{\beta}\mathbf{t}} \left( \left( \mathbf{I} - \hat{\mathcal{F}}(\lambda_{\mathbf{n}}^{\beta})(\lambda_{\mathbf{n}}^{\beta}) \right) \boldsymbol{\Psi}_{\mathbf{0}}, \boldsymbol{\Phi}_{\mathbf{n}}^{\beta*} \right) \boldsymbol{\Phi}_{\mathbf{n}}^{\beta} + \\ & \sum_{\mathbf{n} \in \mathbb{Z}'} \mathbf{e}^{\lambda_{\mathbf{n}}^{\delta}\mathbf{t}} \left( \left( \mathbf{I} - \hat{\mathcal{F}}(\lambda_{\mathbf{n}}^{\delta})(\lambda_{\mathbf{n}}^{\delta}) \right) \boldsymbol{\Psi}_{\mathbf{0}}, \boldsymbol{\Phi}_{\mathbf{n}}^{\delta*} \right) \boldsymbol{\Phi}_{\mathbf{n}}^{\delta} + \\ & \frac{1}{\pi} \int_{\mathbf{0}}^{\infty} \mathbf{e}^{-\mathbf{r}\mathbf{t}} \left( \Im \mathcal{R}(-\mathbf{r}) \left( \mathbf{I} - \hat{\mathcal{F}}(-\mathbf{r}) \right) \boldsymbol{\Psi}_{\mathbf{0}} \right) \mathbf{d}\mathbf{r} \end{split}$$



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Main model equation:  $\dot{\Psi}(t)=i\left(\mathcal{L}_{\beta\delta}-i\mathfrak{M}\right)\Psi(t)+\int\limits_{0}^{t}G(t-\tau)\Psi(\tau)d\tau$ 

-Reduced model equation:  $\dot{\Psi}(t)=i\mathcal{K}_{eta\delta}\Psi(t),\quad \mathcal{K}_{eta\delta}=\mathcal{L}_{eta\delta}-i\mathfrak{M}.$ 

$$\widetilde{\mathbf{S}}\ddot{\mathbf{h}} + \widetilde{\mathbf{S}}\ddot{\alpha} + \mathbf{E}\mathbf{h}'''' - \pi\rho\mathbf{u}\dot{\mathbf{h}} + \pi\rho\mathbf{u}\left(\frac{3}{2} - \mathbf{a}\right)\dot{\alpha} + \pi\rho\mathbf{u}^{2}\alpha = \mathbf{0}$$

$$\widetilde{\mathbf{S}}\ddot{\mathbf{h}} + \widetilde{\mathbf{I}}\ddot{\alpha} - \mathbf{G}\alpha'' - \pi\rho\mathbf{u}\left(\frac{1}{2} + \mathbf{a}\right)\dot{\mathbf{h}} + \pi\rho\mathbf{u}\left(\frac{1}{2} - \mathbf{a}\right)^{2}\dot{\alpha} - \pi\rho\mathbf{u}^{2}\left(\frac{1}{2} + \mathbf{a}\right)\alpha = \mathbf{0}$$

Heuristic derivation of the energy functional accounting air-structure interaction;

Non-local in time (or memory -type) functional



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$$\begin{split} \mathbf{E}(\mathbf{t}) &= & \frac{1}{2} \int_{-\mathbf{L}}^{0} \left\{ \mathbf{E} |\mathbf{h}''|^2 + \mathbf{G} |\alpha'|^2 + \widetilde{\mathbf{m}} |\dot{\mathbf{h}}|^2 + \widetilde{\mathbf{I}} |\dot{\alpha}|^2 - \pi \rho \mathbf{u}^2 \left( \frac{1}{2} + \mathbf{a} \right) |\alpha|^2 + \\ & \widetilde{\mathbf{S}} \left[ \dot{\alpha} \dot{\overline{\mathbf{h}}} + \dot{\overline{\alpha}} \dot{\mathbf{h}} \right] + 2\pi \rho \mathbf{u} \int_{0}^{\mathbf{t}} |\dot{\mathbf{h}} + \left( \frac{1}{2} - \mathbf{a} \right) \dot{\alpha} + \frac{\mathbf{u}}{2} \alpha |^2 d\tau d\mathbf{x} \end{split}$$

Here,

- $\frac{1}{2}\int\limits_{-L}^{0}E|h^{\prime\prime}(x,t)|^{2}dx$  Potential energy due to bending displacement
- ullet  $\frac{1}{2}\int\limits_{-L}^{0}\widetilde{m}|\dot{h}(x,t)|^{2}dx$  Kinetic energy due to bending displacement
- $\frac{1}{2}\int\limits_{-L}^{0}G|\alpha'(x,t)|^2dx$  Potential energy due to torsional displacement



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$$\begin{split} \mathbf{E}(\mathbf{t}) &= & \frac{1}{2} \int\limits_{-\mathbf{L}}^{0} \{ \mathbf{E} |\mathbf{h}''|^2 + \mathbf{G} |\alpha'|^2 + \widetilde{\mathbf{m}} |\dot{\mathbf{h}}|^2 + \widetilde{\mathbf{I}} |\dot{\alpha}|^2 - \pi \rho \mathbf{u}^2 (\frac{1}{2} + \mathbf{a}) |\alpha|^2 + \\ & \widetilde{\mathbf{S}} [\dot{\alpha} \dot{\overline{\mathbf{h}}} + \dot{\overline{\alpha}} \dot{\mathbf{h}}] + 2\pi \rho \mathbf{u} \int\limits_{0}^{t} |\dot{\mathbf{h}} + \left(\frac{1}{2} - \mathbf{a}\right) \dot{\alpha} + \frac{\mathbf{u}}{2} \alpha |^2 \} \mathbf{d}\tau \} \mathbf{d}\mathbf{x} \end{split}$$

Here,

- $\frac{1}{2}\int\limits_{-L}^{0}\widetilde{I}|\dot{\alpha}(x,t)|^{2}dx$  Kinetic energy due to torsional displacement
- $\frac{1}{2}\int\limits_{-L}^{0}\widetilde{S}[\dot{\alpha}\dot{\overline{h}}+\dot{\overline{\alpha}}\dot{h}](x,t)dx$  Kinetic energy due to bending-torsion coupling
- $\bullet \qquad \pi \rho u \int\limits_{-L}^{0} dx \int\limits_{0}^{t} |\left(\dot{h} + \left(\frac{1}{2} a\right) \dot{\alpha} + \frac{u}{2} \alpha\right) (x,\tau)|^{2} d\tau \text{ Energy of vibration due to air-structure interaction}$



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Theorem. For each value of the airspeed u, there exists  $\mathcal{R}(u) > 0$  such that the following statement holds. If an eigenvalue  $\lambda_n$  satisfies  $|\lambda_n| > \mathcal{R}(u)$ , then this eigenmode is stable, i. e.,  $\Re \lambda_n < 0$ . The following estimate is valid for  $\mathcal{R}(u)$ :

$$\mathcal{R}(u) = \mathcal{C}\sqrt{\frac{\rho}{G \Re \delta}} \ u^{3/2}$$

with C being an absolute constant.

Corollary. For each u, all unstable modes are located inside the "circle of instability"  $|\lambda| = \mathcal{R}(u)$ . The number of these eigenmodes is finite.



# Possio Integral Equation in Aeroelasticity

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Physical assumptions on airflow

- 1. Flow  $\Rightarrow$  nonviscous fluid  $\Rightarrow$  Euler equation for stream velocity  $\vec{v}$
- 2. Compressible Flow:  $\rho_t + \nabla \bullet (\rho \vec{v}) = 0$  (continuity equation)
- 3. Isentropic flow:  $P = k\rho^{\gamma}$
- 4. Irrotational (or potential) flow:  $\vec{v} = \nabla \Phi$
- 5. Subsonic (0 < M < 1)
- 6. Coupling between structure and airflow
  - Flow Tangency Condition
     (Flow is attached to wing surface)
  - Kutta-Joukowsky Condition (Pressure drop off the wing and on trailing edge)



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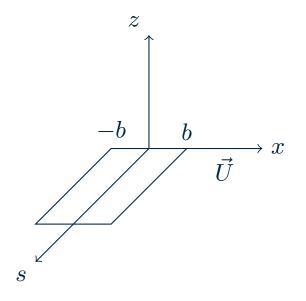
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The Possio Integral Equation relates pressure distribution over a typical section of a slender wing to a normal velocity of the points on a wing surface ("downwash")



Model: High aspect-ratio planar wing; all cross-sections along wing-span are identical; only one spatial variable along cord

• Velocity field  $\to$  potential U-free stream velocity Velocity potential:

$$U + \phi(x, z, t),$$

$$-\infty < x < \infty, \quad 0 < z < \infty,$$

$$t > 0$$



# Linearized version of Euler Equation for disturbance potential

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$$\frac{\partial^2 \phi}{\partial^2 t} + 2M a_{\infty} \frac{\partial^2 \phi}{\partial t \partial x} = a_{\infty}^2 (1 - M^2) \frac{\partial^2 \phi}{\partial x^2} + a_{\infty}^2 \frac{\partial^2 \phi}{\partial x^2}$$

 $a_{\infty}$ -speed of sound in flight altitude  $M=U/a_{\infty}$ -Mach number

Boundary conditions make the problem complicated

- Flow Tangency Condition (flow is attached)
- Kutta-Joukowsky Condition (zero pressure off the wing and at the trailing edge)
- Far Field conditions



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1. Flow Tangency Condition

$$\frac{\partial}{\partial z}\phi(x,z,t)|_{z=0} = w_a(x,t), |x| < b$$

 $w_a$ -given normal velocity of wing

2. Kutta-Joukowsky Condition Acceleration potential:

$$\psi(x, z, t) = \frac{\partial \phi}{\partial t} + U \frac{\partial \phi}{\partial x}$$

- $\psi(x, 0, t) = 0, |x| > b$
- $\bullet \quad \lim_{x \to b-0} \psi(x,0,t) = 0$
- 3. Far Field conditions 
  Disturbance potential o 0 as  $|x| o \infty$  or  $z o \infty$



# The Possio Integral Equation

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$$W_{a}(\cdot,\lambda) = \frac{\sqrt{1-M^{2}}}{2} \left\{ \mathcal{H}_{b}A(\cdot,\lambda) - \frac{1}{\sqrt{1-M^{2}}} \times \left[ \tilde{\lambda}\mathcal{H}_{b}\mathcal{L}(\tilde{\lambda})A(\cdot,\lambda) - \tilde{\lambda}g_{-}(\tilde{\lambda},x)e^{-b\tilde{\lambda}}L\left(\tilde{\lambda},A(\cdot,\lambda)\right) \right] - \int_{0}^{\alpha_{1}} a(s) \left[ \tilde{\lambda}\mathcal{H}_{b}\mathcal{L}(\tilde{\lambda}s)A(\cdot,\lambda) - \tilde{\lambda}g_{-}(\tilde{\lambda}s,x)e^{-\tilde{\lambda}bs}L\left(\tilde{\lambda}s,A(\cdot,\lambda)\right) \right] ds + \int_{0}^{\alpha_{2}} a(-s) \left[ \tilde{\lambda}\mathcal{H}_{b}\mathcal{L}^{*}(\tilde{\lambda}s)A(\cdot,\lambda) + \tilde{\lambda}g_{+}(\tilde{\lambda}s,x)e^{-\tilde{\lambda}bs}L\left(-\tilde{\lambda}s,A(\cdot,\lambda)\right) \right] ds \right\}$$



# Asymptotical form of Possio Equation as $\lambda \to \infty$

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$$2\mathcal{T}W_{a}(\cdot,\lambda) = \left[\sqrt{1-M^{2}} - \tilde{\lambda}\mathcal{L}(\tilde{\lambda})\right]F(\cdot,\lambda) + \\ \tilde{\lambda}e^{-b\tilde{\lambda}}h_{-}(x,\tilde{\lambda})L(\tilde{\lambda},F(\cdot,\lambda)) - \\ M\left\{\mathcal{H}_{b}[F(\cdot,\lambda)] - \mathcal{T}[F(\cdot,\lambda)]\right\}$$

# Main difficulty:

 $\mathcal{L}(\lambda)$   $\rightarrow$  Volterra integral operator

 $L(\tilde{\lambda},\cdot)$   $\to$  Integral operator with degenerate kernel

 $\mathcal{H}_b(\cdot) \quad o \quad$  Finite Hilbert transform

 $\mathcal{T}(\cdot)$  ightarrow "Inverse" to  $\mathcal{H}_b$ 

 $\mathcal{H}_b$  and  $\mathcal{T} o$  Singular integral operators